UDC 338.984 Articles

doi:10.31799/1684-8853-2018-6-14-23

# Exact solution method for Fredholm integro-differential equations with multipoint and integral boundary conditions. Part 1. Extention method

**N. N. Vassiliev<sup>a,b</sup>,** PhD, Tech., Senior Researcher, orcid.org/0000-0002-0841-1168, vasiliev@pdmi.ras.ru **I. N. Parasidis<sup>c</sup>,** PhD, Associate Professor, paras@teilar.gr

E. Providas<sup>d</sup>, PhD, Associate Professor, providas@teilar.gr

<sup>a</sup>Saint-Petersburg Department of V. A. Steklov Institute of Mathematics of the RAS, 27, Fontanka, 191023, Saint-Petersburg, Russian Federation

<sup>b</sup>Saint-Petersburg Electrotechnical University ETU "LETI", 5, Professora Popova St., 197376, Saint-Petersburg, Russian Federation

<sup>c</sup>Department of Electrical Engineering, Technological Educational Institute of Thessaly, 41110, Larissa, Greece

<sup>d</sup>Department of Mechanical Engineering, Technological Educational Institute of Thessaly, 41110, Larissa, Greece

**Introduction:** Boundary value problems for differential and integro-differential equations with multipoint and non-local boundary conditions often arise in mechanics, physics, biology, biotechnology, chemical engineering, medical science, finances and other fields. Finding an exact solution of a boundary value problem with Fredholm integro-differential equations is a challenging problem. In most cases, solutions are obtained by numerical methods. **Purpose:** Search for necessary and sufficient solvability conditions for abstract operator equations and their exact solutions. **Results:** A direct method is proposed for the exact solution of a certain class of ordinary differential or Fredholm integro-differential equations with separable kernels and multipoint/integral boundary conditions. We study abstract equations of the form Bu = Au - gF(Au) = f and  $B_1u = A^2u - gF(Au) - gF(A^2u) = f$  with non-local boundary conditions  $\Phi(u) = N\Psi(u)$  and  $\Phi(u) = N\Psi(u)$ ,  $\Phi(Au) = DF(Au) + N\Psi(Au)$ , respectively, where A is a differential operator, g and g are vectors, g and g are reactors, g and g are reactors, g and g are matrices, and g are functional vectors. This method is simple to use and can be easily incorporated into any Computer Algebra System (CAS). The upcoming Part 2 of this paper will be devoted to decomposition method for this problem where the operator g is quadratic factorable.

**Keywords** — differential and Fredholm integro-differential equations, multipoint and non-local integral boundary conditions, correct operators, exact solutions.

Citation: Vassiliev N. N., Parasidis I. N., Providas E. Exact solution method for Fredholm integro-differential equations with multipoint and integral boundary conditions. Part 1. Extension method. *Informatsionno-upravliaiushchie sistemy* [Information and Control Systems], 2018, no. 6, pp. 14–23. doi:10.31799/1684-8853-2018-6-14-23

# Introduction

Boundary value problems (BVP) for differential and integro-differential equations (IDE) with multipoint and nonlocal boundary conditions arise in various fields of mechanics, physics, biology, biotechnology, chemical engineering, medical science, finance and others [1-14]. More precisely these are elasticity, heat and mass transfer, diffraction, underground water flow and population dynamics problems. Perhaps the first known problem which was reduced to the IDE  $a_1 y^{iv}(t) + y(t) = -a_2 \int_{-1}^{1} K(t, x) y^{iv}(x) dx$  is Proctor's problem of Equilibrium of an elastic beam in XVII century. Fredholm integro-differential equations with nonlocal integral boundary conditions and ordinary differential operators, probably, first were considered by J. D. Tamarkin [15]. Problems with nonlocal boundary conditions for elliptic equations first were investigated by A. V. Bitsadze, A. A. Samarskii [16], while BVP for parabolic equations with nonlocal integral boundary conditions were studied by J. R. Cannon [5], L. I. Kamynin [7], N. I. Ionkin [6] and others. Later such investigations for Laplace, Poisson and heat equations were explored by V. A. Il'in and E. L. Moiseev [17] and others [18-20]. Nonlocal BVP involving integral conditions for hyperbolic equations were studied in [21]. Multipoint and nonlocal BVP with integral boundary conditions for ordinary differential equations were considered in [22, 23]. Fractional IDE with integral boundary conditions were given in [24]. The problem of the existence of solutions for nonlocal BVP was the subject of many papers [19, 20, 23, 25-28]. Exact solutions of BVP with Fredholm IDE were considered in [29] and [30]. In most cases numerical methods are employed. Here, the necessary and sufficient solvability conditions of the abstract operator equations:

$$Bu = Au - Qu, Qu = gF(Au),$$

$$\mathfrak{D}(B) = \{u \in \mathfrak{D}(A) : \Phi(u) = N\Psi(u)\};$$
(1)

$$B_1 u = A^2 u - Q_1 u, Q_1 u = \mathbf{q} \mathbf{F}(A u) + \mathbf{g} \mathbf{F}(A^2 u),$$

$$\mathfrak{D}(B_1) = \{ u \in \mathfrak{D}(A^2) : \mathbf{\Phi}(u) = \mathbf{N} \mathbf{\Psi}(u) \},$$

$$\mathbf{\Phi}(A u) = \mathbf{D} \mathbf{F}(A u) + \mathbf{N} \mathbf{\Psi}(A u),$$
(2)

and their exact solutions are obtained in closed form. This formalism is applied to solve Fredholm IDE with multipoint or nonlocal integral boundary conditions, when A is a differential operator and Q,  $Q_1$  are integral operators with separable kernels. The problems (1), (2) arise naturally from A. A. Dezin, R. O. Oinarov extensions of linear operators [31, 26], which are not restrictions of a maximal operator, unlike the classical M. G. Krein, J. Von. Neuman extensions [32, 33] in Hilbert space and in Banach space [34]. This work is a generalization of the papers [26-28, 35], where integral boundary conditions have not been considered. Solving differential or Fredholm IDE with integral boundary conditions is a complicated problem, since the operators B and  $B_1$  in (1), (2) are obtained by perturbations of boundary conditions and the action of an operator A. Whereas in [26–28, 35] the operators  $B = \hat{A} + Q$ ,  $\mathfrak{D}(B) = \mathfrak{D}(\hat{A})$  and  $B_1 = \hat{A}^2 + Q_1$ ,  $\mathfrak{D}(B_1) = \mathfrak{D}(\hat{A}^2)$ are obtained only by perturbation of the action of a correct operator  $\hat{A}$  which is a restriction of a maximal operator A.

### Terminology and notation

Let X, Y be complex Banach spaces and  $X^*$  the adjoint space of X, i. e. the set of all complex-valued linear and bounded functionals on X. We denote by f(x) the value of f on x. We write  $\mathfrak{D}(A)$  and R(A) for the domain and the range of the operator A, respectively. An operator  $A_2$  is said to be an extension of an operator  $A_1$ , or  $A_1$  is said to be a restriction of  $A_2$ , in symbol  $A_1 \subset A_2$ , if  $\mathfrak{D}(A_2) \supseteq \mathfrak{D}(A_1)$  and  $A_1x = A_2x$ , for all  $x \in \mathfrak{D}(A_1)$ . An operator  $A: X \to Y$  is called closed if for every sequence  $x_n$  in  $\mathfrak{D}(A)$  converging to  $x_0$  with  $Ax_n \to f_0$ , it follows that  $x_0 \in \mathfrak{D}(A)$  and  $Ax_0 = f_0$ . A closed operator A is called maximal if R(A) = Y and ker  $A \neq \{0\}$ . An operator  $\hat{A}: X \to Y$  is called correct if  $R(\hat{A}) = Y$  and the inverse  $\hat{A}^{-1}$ exists and is continuous on Y. An operator  $\hat{A}$  is called a correct restriction of the maximal operator A if it is a correct operator and  $\hat{A} \subset A$ . If  $\Psi_i \in X^*$ , i=1, ..., n, then we denote by  $\Psi = col(\Psi_1, ..., \Psi_n)$  and  $\Psi(x) = col(\Psi_1(x), ..., \Psi_n(x))$ . Let  $g = (g_1, ..., g_n)$  be a vector of  $X^n$ . We will denote by  $\Psi(g)$  the  $n \times n$  matrix whose i, j-th entry  $\Psi_i(g_i)$  is the value of functional  $\Psi_i$  on element  $g_i$ . Note that  $\Psi(gC) = \Psi(g)C$ , where C is a  $n \times k$  constant matrix. We will also denote by  $\mathbf{0}_n$  the zero and by  $\mathbf{I}_n$  the identity  $n \times n$ matrices. By 0 we will denote the zero column vector.

# Extension methods for ordinary differential and Fredholm IDE

Let  $A: X \xrightarrow{on} X$  be an ordinary  $m^{\text{th}}$  order differential operator

$$Au(x) = \alpha_0 u^{(m)}(x) + \alpha_1 u^{(m-1)}(x) + \dots + \alpha_m u(x),$$
 
$$\alpha_i \in \mathbb{R}$$
 (3)

and X be a Banach space. Usually  $X=\mathbb{C}[a,\ b]$  or  $X=L_p(a,\ b)$ ,  $p\geq 1$ . In the sequel we denote by  $X_A^m=\left(D(A),\ \|\cdot\|_{X^m}\right)$  the Banach space of all m times differentiable functions with norm  $\|u(x)\|_{X_A^m}=\sum_{i=0}^m \left\|u^{(i)}(x)\right\|_X$  and by  $X_A^{m-1}$  the Banach space of all

$$m-1$$
 times differentiable functions with norm

$$||u(x)||_{X_A^{m-1}} = \sum_{i=0}^{m-1} ||u^{(i)}(x)||_X.$$
 (4)

Note that for  $X=\mathbb{C}[a,\,b]$  the spaces  $X_A^m$ ,  $X_A^{m-1}$  are defined by  $\mathbb{C}^m[a,\,b]$ ,  $\mathbb{C}^{m-1}[a,\,b]$ , respectively. It is a well-known fact that the operator defined by

$$\hat{A}u(x) = \alpha_0 u^{(m)}(x) +$$

$$+\alpha_1 u^{(m-1)}(x) + \dots + \alpha_m u(x) = f,$$

$$\alpha_i \in \mathbb{R}, x \in [a, b],$$

$$\mathfrak{D}(\hat{A}) =$$

$$(5)$$

$$= \left\{ u(x) \in \mathbb{C}^m [a, b] : u(a) = u'(a) = \dots = u^{m-1}(a) = 0 \right\}$$

is a correct restriction of A and the unique solution of (5) is

$$u(x) = \hat{A}^{-1}f(x) = \frac{1}{(m-1)!} \int_{a}^{x} (x-t)^{m-1} f(t) dt,$$
$$f(x) \in \mathbb{C}[a,b]. \tag{6}$$

**Lemma 1.** Let  $A_i$ ,  $B_i$ ,  $C_i$ , D are  $n \times n$  matrices,

where 
$$i = 1, 2, 3$$
, and  $G = \begin{pmatrix} A_1 & A_2 & A_3 \\ B_1 & B_2 & B_3 \\ C_1 & C_2 & C_3 \end{pmatrix}$ . Then the

next properties of determinants hold true:

$$\det \begin{pmatrix} \mathbf{A}_{1} & \mathbf{A}_{2} & \mathbf{A}_{3} \\ \mathbf{B}_{1} & \mathbf{B}_{2} & \mathbf{B}_{3} \\ \mathbf{C}_{1} & \mathbf{C}_{2} & \mathbf{C}_{3} \end{pmatrix} = \\ = \det \begin{pmatrix} \mathbf{A}_{1} \pm \mathbf{D}\mathbf{B}_{1} & \mathbf{A}_{2} \pm \mathbf{D}\mathbf{B}_{2} & \mathbf{A}_{3} \pm \mathbf{D}\mathbf{B}_{3} \\ \mathbf{B}_{1} & \mathbf{B}_{2} & \mathbf{B}_{3} \\ \mathbf{C}_{1} & \mathbf{C}_{2} & \mathbf{C}_{3} \end{pmatrix}; \qquad (7)$$

### ΤΕΟΡΕΤИЧЕСКАЯ И ΠΡИΚΛΑΔΗΑЯ ΜΑΤΕΜΑΤИΚΑ

$$\det \begin{pmatrix} \mathbf{A}_{1} & \mathbf{A}_{2} & \mathbf{A}_{3} \\ \mathbf{B}_{1} & \mathbf{B}_{2} & \mathbf{B}_{3} \\ \mathbf{C}_{1} & \mathbf{C}_{2} & \mathbf{C}_{3} \end{pmatrix} = \det \begin{pmatrix} \mathbf{A}_{1} & \mathbf{A}_{2} \pm \mathbf{A}_{3} \mathbf{D} & \mathbf{A}_{3} \\ \mathbf{B}_{1} & \mathbf{B}_{2} \pm \mathbf{B}_{3} \mathbf{D} & \mathbf{B}_{3} \\ \mathbf{C}_{1} & \mathbf{C}_{2} \pm \mathbf{C}_{3} \mathbf{D} & \mathbf{C}_{3} \end{pmatrix}. \quad (8)$$

$$Proof: \quad \text{Let} \quad \mathbf{H} = \begin{pmatrix} \mathbf{I}_n & -\mathbf{D} & \mathbf{0}_n \\ \mathbf{0}_n & \mathbf{I}_n & \mathbf{0}_n \\ \mathbf{0}_n & \mathbf{0}_n & \mathbf{I}_n \end{pmatrix}. \quad \text{Then} \quad \mathbf{H}^{-1} =$$

$$= \begin{pmatrix} \mathbf{I}_n & \mathbf{D} & \mathbf{0}_n \\ \mathbf{0}_n & \mathbf{I}_n & \mathbf{0}_n \\ \mathbf{0}_n & \mathbf{0}_n & \mathbf{I}_n \end{pmatrix}, \quad \det \mathbf{H} = \det \mathbf{H}^{-1} = \mathbf{1}, \quad |\mathbf{H}\mathbf{G}| = |\mathbf{H}||\mathbf{G}| =$$

= |G| and  $|H^{-1}G| = |H^{-1}||G| = |G|$ . So (7) holds.

$$= \begin{pmatrix} \mathbf{I}_n & \mathbf{0}_n & \mathbf{0}_n \\ \mathbf{0}_n & \mathbf{I}_n & \mathbf{0}_n \\ \mathbf{0}_n & -\mathbf{D} & \mathbf{I}_n \end{pmatrix}, \quad |\mathbf{H}| = |\mathbf{H}^{-1}| = 1, \quad |\mathbf{G}\mathbf{H}| = |\mathbf{G}||\mathbf{H}| = |\mathbf{G}|$$

and  $|GH^{-1}| = |G||H^{-1}| = |G|$ . So (8) holds and Lemma 1 is proved.

matrices. Let  $\Gamma$  be the matrix obtained from G by multiplying from the left a row by the  $n \times n$  matrix D and then adding it to another row, or by multiplying from the right a column of G by the matrix D and then adding it to another column of G. Then  $\det G = \det \Gamma$ .

**Theorem 1.** Let X be a complex Banach space,  $a: X \to X$  an operator from (3) with finite dimensional kernel  $\mathbf{z} = (z_1, ..., z_m)$  which is a basis of ker A, and let  $\hat{A}$  be a correct restriction of A defined by

$$\hat{A} \subset A,$$

$$\mathfrak{D}(\hat{A}) = \{ u \in \mathfrak{D}(A) : \Phi(u) = \mathbf{0} \},$$
(9)

the components of the functional vectors  $\mathbf{\Phi} = col\ (\mathbf{\Phi}_1,\ ...,\ \mathbf{\Phi}_m),\ \mathbf{\Psi} = col\ (\mathbf{\Psi}_1,\ ...,\ \mathbf{\Psi}_n)$  and  $\mathbf{F} = col\ (F_1,\ ...,\ F_n)$  belong to  $X^{m-1}$  and respectively.

Suppose also that  $\Phi_1$ , ...,  $\Phi_m$  biorthogonal to  $z_1$ , ...,  $z_m$  and that the components of vector  $\mathbf{g} = (g_1, ..., g_n) \in X^n$  are linearly independent and  $\mathbf{N}$  is a  $m \times n$  matrix. Then:

(i) The operator *B* defined by

$$Bu = Au - \mathbf{g}\mathbf{F}(Au) = f,$$
 
$$f \in X;$$
 
$$\mathfrak{D}(B) = \{u \in \mathfrak{D}(A) : \mathbf{\Phi}(u) = \mathbf{N}\mathbf{\Psi}(u)\}$$
 (10)

is injective if and only if

$$\begin{aligned} \det &\mathbf{V} = \det[\mathbf{I}_n - \Psi(\mathbf{z})\mathbf{N}] \neq 0 \text{ and} \\ &\det &\mathbf{W} = \det[\mathbf{I}_n - \mathbf{F}(\mathbf{g})] \neq 0. \end{aligned} \tag{11}$$

(ii) If B is injective, then B is correct and for all  $f \in X$  the unique solution of (10) is given by

$$u = B^{-1}f = \hat{A}^{-1}f + \left[\hat{A}^{-1}\mathbf{g} + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\mathbf{\Psi}(\hat{A}^{-1}\mathbf{g})\right] \times \mathbf{W}^{-1}\mathbf{F}(f) + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\mathbf{\Psi}(\hat{A}^{-1}f).$$
(12)

Proof: (i). Let  $\det \mathbf{W} \neq \mathbf{0}$ ,  $\det \mathbf{V} \neq \mathbf{0}$  and  $u \in \ker B$ . Then  $Bu = Au - \mathbf{gF}(Au) = \mathbf{0}$ ,  $\Phi(u) = \mathbf{N}\Psi(u)$  and  $[\mathbf{I}_n - \mathbf{F}(\mathbf{g})]\mathbf{F}(Au) = \mathbf{0}$ ,  $\Phi(u - \mathbf{z}\mathbf{N}\Psi(u)) = \mathbf{0}$ . The last equation, since (9), implies  $u - \mathbf{z}\mathbf{N}\Psi(u) \in \mathfrak{D}(\hat{A})$ . From  $[\mathbf{I}_n - \mathbf{F}(\mathbf{g})]\mathbf{F}(Au) = \mathbf{0}$ , since  $\det W \neq \mathbf{0}$ , follows  $\mathbf{F}(Au) = \mathbf{0}$ . Then  $Bu = Au = \mathbf{0}$  which yelds  $\hat{A}(u - \mathbf{z}\mathbf{N}\Psi(u)) = \mathbf{0}$  and so  $u = \mathbf{z}\mathbf{N}\Psi(u)$ . Then  $\Psi(u) = \Psi(\mathbf{z})\mathbf{N}\Psi(u)$  or  $[\mathbf{I}_n - \Psi(\mathbf{z})\mathbf{N}]\Psi(u) = \mathbf{0}$ . The last, since  $\det \mathbf{V} \neq \mathbf{0}$  implies  $\Psi(u) = \mathbf{0}$  and so from  $u = \mathbf{z}\mathbf{N}\Psi(u)$  we get  $u = \mathbf{0}$ , i. e.  $\ker B = \{\mathbf{0}\}$  and B is an injective operator.

Conversely. Let  $\det \mathbf{V} = \mathbf{0}$ . Then there exists a vector  $\mathbf{c} = col(c_1, ..., c_n) = \mathbf{0}$  such that  $\mathbf{V}\mathbf{c} = \mathbf{0}$ .

Consider the element  $u_0 = \mathbf{z} \mathbf{Nc} \neq \mathbf{0}$ , otherwise  $\mathbf{Nc} = \mathbf{0}$  and from  $[\mathbf{I}_n - \Psi(\mathbf{z})\mathbf{N}]\mathbf{c} = \mathbf{0}$  follows  $\mathbf{c} = \mathbf{0}$ , which contradicts the hypothesis  $\mathbf{c} \neq \mathbf{0}$ . Note that  $u_0 \in \mathfrak{D}(B)$ , since  $\Phi(u_0) = \mathbf{Nc}$ ,  $\Psi(u_0) = \Psi(\mathbf{z})\mathbf{Nc}$ ,  $\Phi(u_0) - \mathbf{N}\Psi(u_0) = \mathbf{Nc} - \mathbf{N}\Psi(\mathbf{z})\mathbf{Nc} = \mathbf{N}[\mathbf{I}_n - \Psi(\mathbf{z})\mathbf{N}]\mathbf{c} = \mathbf{N}\mathbf{Vc} = \mathbf{0}$ . It is evident that  $u_0 \in \ker B$ . So  $u_0 \in \ker B$ . Hence  $\ker B \neq \{0\}$  and B is not injective. Let now  $\det \mathbf{V} \neq \mathbf{0}$ , but  $\det \mathbf{W} = \mathbf{0}$ . Then there exists a vector  $\mathbf{c} = \operatorname{col}(c_1, \ldots, c_n) \neq \mathbf{0}$  such that  $\mathbf{Wc} = \mathbf{0}$ . Note that  $\mathbf{gc} \neq \mathbf{0}$  because of  $g_1, \ldots, g_n$  is a linearly independent set and that the element  $u_0 = \left[\hat{A}^{-1}\mathbf{g} + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\Psi(A^{-1}\mathbf{g})\right]\mathbf{c} \neq \mathbf{0}$ , otherwise  $\mathbf{g} = \mathbf{0}$ . For  $u_0$  we obtain

$$\begin{split} u_0 = & \left[ \hat{A}^{-1}\mathbf{g} + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\boldsymbol{\Psi} \Big( A^{-1}\mathbf{g} \Big) \right] \mathbf{c} \neq \mathbf{0}, \\ \boldsymbol{\Phi} \Big( u_0 \Big) - \mathbf{N}\boldsymbol{\Psi} \Big( u_0 \Big) &= \mathbf{N}\mathbf{V}^{-1}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} - \mathbf{N}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} - \\ &- \mathbf{N}\boldsymbol{\Psi} \Big( \mathbf{z} \Big) \mathbf{N}\mathbf{V}^{-1}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} = \\ &= \mathbf{N} \Big[ \mathbf{I}_n - \boldsymbol{\Psi} \Big( \mathbf{z} \Big) \mathbf{N} \Big] \mathbf{V}^{-1}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} - \mathbf{N}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} = \\ &= \mathbf{N}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} - \mathbf{N}\boldsymbol{\Psi} \Big( \hat{A}^{-1}\mathbf{g} \Big) \mathbf{c} = \mathbf{0}, \\ & Bu_0 = Au_0 - \mathbf{g}\mathbf{F} \Big( \hat{A}u_0 \Big) = \\ &= \mathbf{g}\mathbf{c} - \mathbf{g}\mathbf{F} \Big( \mathbf{g} \Big) \mathbf{c} = \mathbf{g} \Big[ \mathbf{I}_n - \mathbf{F} \Big( \mathbf{g} \Big) \Big] \mathbf{c} = \mathbf{g}\mathbf{W}\mathbf{c} = \mathbf{g}\mathbf{0} = \mathbf{0}. \end{split}$$

So  $u_0 \in \ker B$ . Consequently  $\ker B \neq \{0\}$  and B is not injective. Hence B is injective if and only if  $\det \mathbf{V} \neq \mathbf{0}$ ,  $\det \mathbf{W} \neq \mathbf{0}$ . The statement (i) holds.

(ii) Let  $\det \mathbf{W} \neq \mathbf{0}$  and  $\det \mathbf{V} \neq \mathbf{0}$ . By statement (i), the operator B is injective. Since  $\mathbf{z} \in [\ker A]^m$ ,  $\Phi(\mathbf{z}) = \mathbf{I}_m$ , the problem (10) is written as

$$Bu = A(u - \mathbf{z}\mathbf{N}\Psi(u)) - \mathbf{g}\mathbf{F}(Au) = f,$$
  

$$f \in X;$$
  

$$\mathfrak{D}(B) = \{u \in \mathfrak{D}(A) : \Phi(u - \mathbf{z}\mathbf{N}\Psi(u)) = 0\}.$$
 (13)

Then, applying Equation (9) and relation  $Bu = \hat{A}(u - \mathbf{z}\mathbf{N}\Psi(u)) - \mathbf{g}\mathbf{F}(Au) = f$  $u - \mathbf{z} \mathbf{N} \Psi(u) \in \mathfrak{D}(\hat{A}), Bu = \hat{A}(u - \mathbf{z} \mathbf{N} \Psi(u)) - \mathbf{g} \mathbf{F}(Au) = f$ and for every  $u \in \mathfrak{D}(B)$ ,  $f \in X$  using (10), (13) we obtain

$$\begin{split} & \left[\mathbf{I}_{n} - \mathbf{F}(\mathbf{g})\right] \mathbf{F}(Au) = \mathbf{F}(f), \\ & \mathbf{F}(Au) = \mathbf{W}^{-1} \mathbf{F}(f), \\ & u - \mathbf{z} \mathbf{N} \Psi(u) = \hat{A}^{-1} \mathbf{g} \mathbf{F}(Au) + \hat{A}^{-1} f, \\ & \Psi(u) - \Psi(\mathbf{z}) \mathbf{N} \Psi(u) = \Psi(\hat{A}^{-1} \mathbf{g}) \mathbf{F}(Au) + \Psi(\hat{A}^{-1} f), \\ & \left[\mathbf{I}_{n} - \Psi(\mathbf{z}) \mathbf{N}\right] \Psi(u) = \Psi(\hat{A}^{-1} \mathbf{g}) \mathbf{W}^{-1} \mathbf{F}(f) + \Psi(\hat{A}^{-1} f), \\ & \Psi(u) = \mathbf{V}^{-1} \left[\Psi(\hat{A}^{-1} \mathbf{g}) \mathbf{W}^{-1} \mathbf{F}(f) + \Psi(\hat{A}^{-1} f)\right], \\ & u = B^{-1} f = \hat{A}^{-1} f + \hat{A}^{-1} \mathbf{g} \mathbf{W}^{-1} \mathbf{F}(f) + \\ & + \mathbf{z} \mathbf{N} \mathbf{V}^{-1} \left[\Psi(\hat{A}^{-1} \mathbf{g}) \mathbf{W}^{-1} \mathbf{F}(f) + \Psi(\hat{A}^{-1} f)\right]. \end{split}$$

From the last equation for every  $f \in X$  follows the unique solution (12) of (10). Because f in (12) is arbitrary, we obtain R(B) = X. Since the operator  $\hat{A}^{-1}$  and the functionals  $F_1, ..., F_n, \Psi_1, ..., \Psi_n$  are bounded, from (12) follows the boundedness of  $B^{-1}$ . Hence, the operator B is correct if and only if (11) holds and the unique solution of (10) is given by (12). The theorem is proved.

From the previous theorem for g = 0 follows the next corollary which is useful for solving some classes of differential equations with nonlocal boundary conditions.

**Corollary 1.** Let a complex Banach space X, the operators A,  $\hat{A}$ , the vector  $\mathbf{z}$  and functional vectors  $\Phi$ ,  $\Psi$  and the matrix **N** be defined as in Theorem 1. Then:

(i) The operator B defined by

$$Bu = Au = f,$$
 
$$f \in X;$$
 
$$\mathfrak{D}(B) = \{ u \in \mathfrak{D}(A) : \Phi(u) = \mathbb{N}\Psi(u) \}$$
 (14)

is correct if and only if  $\det \mathbf{V} = \det [\mathbf{I}_n - \Psi(\mathbf{z})\mathbf{N}] \neq 0$  and for all  $f \in X$  the unique solution of (14) is given by

$$u = B^{-1}f = \hat{A}^{-1}f + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\Psi(\hat{A}^{-1}f).$$
 (15)

**Theorem 2.** Let a Banach space X, the vectors  $\mathbf{z}$ ,  $\mathbf{\Phi}$ ,  $\Psi$ , F, the operators A,  $\hat{A}$  be defined as in Theorem 1 and the operator  $B_1: X \to X$  by

$$B_1 u = A^2 u - \mathbf{q} \mathbf{F}(Au) - \mathbf{g} \mathbf{F}(A^2 u) = f;$$

$$\mathfrak{D}(B_1) = \{ u \in \mathfrak{D}(A^2) : \mathbf{\Phi}(u) = \mathbf{N} \mathbf{\Psi}(u),$$

$$\mathbf{\Phi}(Au) = \mathbf{D} \mathbf{F}(Au) + \mathbf{N} \mathbf{\Psi}(Au) \}.$$
(17)

Suppose also that the vectors q and g are linearly independent,  $q = (q_1, ..., q_n), g = (g_1, ..., g_n) \in X^n$ , and **D**, **N** are  $m \times n$  matrices. Then:

(i) The operator  $B_1$  corresponding to the problem (16), (17) is injective if and only if

$$\det \mathbf{L} =$$

$$= \det \begin{bmatrix} \mathbf{0}_{n} & -\mathbf{F}(\mathbf{z})\mathbf{N} & \mathbf{K}_{1} & -\mathbf{F}(\hat{A}^{-1}\mathbf{g}) \\ \mathbf{V} & -\Psi(\hat{A}^{-1}\mathbf{z})\mathbf{N} & -\mathbf{K}_{3} & -\Psi(\hat{A}^{-2}\mathbf{g}) \\ \mathbf{0}_{n} & \mathbf{V} & -\mathbf{K}_{2} & -\Psi(\hat{A}^{-1}\mathbf{g}) \\ \mathbf{0}_{n} & \mathbf{0}_{n} & -\mathbf{F}(\mathbf{q}) & \mathbf{W} \end{bmatrix} \neq \mathbf{0}, (18)$$
where

$$\begin{split} \mathbf{K}_1 &= \mathbf{I}_n - \mathbf{F}(\mathbf{z}) \mathbf{D} - \mathbf{F}(\hat{A}^{-1}\mathbf{q}), \quad \mathbf{K}_2 &= \Psi(\mathbf{z}) \mathbf{D} + \Psi(\hat{A}^{-1}\mathbf{q}), \\ \mathbf{K}_3 &= \Psi(\hat{A}^{-1}\mathbf{z}) \mathbf{D} + \Psi(\hat{A}^{-2}\mathbf{q}), \\ \mathbf{W} &= \mathbf{I}_n - \mathbf{F}(\mathbf{g}), \quad \mathbf{V} = \mathbf{I}_n - \Psi(\mathbf{z}) \mathbf{N}. \end{split} \tag{19}$$

(ii) If the operator  $B_1$  is injective, then it is correct and the unique solution of (16), (17) is given by

$$u = B_1^{-1} f = \hat{A}^{-2} f +$$

$$+ \left( \mathbf{z} \mathbf{N}, \, \hat{A}^{-1} \mathbf{z} \mathbf{N}, \, \hat{A}^{-1} \mathbf{z} \mathbf{D} + \hat{A}^{-2} \mathbf{q}, \, \hat{A}^{-2} \mathbf{g} \right) \times$$

$$\times \mathbf{L}^{-1} col \left( \mathbf{F} \left( \hat{A}^{-1} f \right), \, \mathbf{\Psi} \left( \hat{A}^{-2} f \right), \, \mathbf{\Psi} \left( \hat{A}^{-1} f \right), \, \mathbf{F} (f) \right).$$
 (20)

*Proof*: (i) Let detL  $\neq$  0. Since  $\Phi(\mathbf{z}) = \mathbf{I}_m$ , the relations (17) can be represented as

$$\Phi(u - \mathbf{z}\mathbf{N}\Psi(u)) = 0,$$

$$\Phi(Au - \mathbf{z}\mathbf{D}\mathbf{F}(Au) - \mathbf{z}\mathbf{N}\Psi(Au)) = 0,$$

which taking into account (9) imply

$$u - \mathbf{z} \mathbf{N} \Psi(u) \in \mathfrak{D}(\hat{A});$$
 (21)

$$Au - \mathbf{z}\mathbf{D}\mathbf{F}(Au) - \mathbf{z}\mathbf{N}\Psi(Au) \in \mathfrak{D}(\hat{A}).$$
 (22)

Then, since  $\mathbf{z} \in [\ker A]^m$ ,  $\hat{A} \subset A$  and  $\hat{A}$  is correct, from (16) we obtain

$$\hat{A} \Big( Au - \mathbf{z} \Big[ \mathbf{DF} (Au) + \mathbf{N\Psi} (Au) \Big] \Big) - \mathbf{qF} (Au) - \mathbf{gF} \Big( A^2 u \Big) = f,$$

### ТЕОРЕТИЧЕСКАЯ И ПРИКЛАДНАЯ МАТЕМАТИКА

$$Au - \mathbf{z} \Big[ \mathbf{DF} (Au) + \mathbf{N\Psi} (Au) \Big] - \hat{A}^{-1} \mathbf{qF} (Au) - \hat{A}^{-1} \mathbf{gF} \Big( A^2 u \Big) = \hat{A}^{-1} f,$$

$$\hat{A} \Big( u - \mathbf{zN\Psi} (u) \Big) - \mathbf{z} \Big[ \mathbf{DF} (Au) + \mathbf{N\Psi} (Au) \Big] - \hat{A}^{-1} \mathbf{qF} (Au) - \hat{A}^{-1} \mathbf{gF} \Big( A^2 u \Big) = \hat{A}^{-1} f,$$

$$u - \mathbf{zN\Psi} (u) - \hat{A}^{-1} \mathbf{z} \Big[ \mathbf{DF} (Au) + \mathbf{N\Psi} (Au) \Big] - \hat{A}^{-2} \mathbf{qF} (Au) - \hat{A}^{-2} \mathbf{gF} \Big( A^2 u \Big) = \hat{A}^{-2} f.$$

Then taking into account (16) we get

$$A^{2}u = \mathbf{q}\mathbf{F}(Au) + \mathbf{g}\mathbf{F}(A^{2}u) + f,$$

$$Au = \mathbf{z}\Big[\mathbf{D}\mathbf{F}(Au) + \mathbf{N}\mathbf{\Psi}(Au)\Big] +$$

$$+\hat{A}^{-1}\mathbf{q}\mathbf{F}(Au) + \hat{A}^{-1}\mathbf{g}\mathbf{F}(A^{2}u) + \hat{A}^{-1}f,$$

$$u = \mathbf{z}\mathbf{N}\mathbf{\Psi}(u) + \hat{A}^{-1}\mathbf{z}\mathbf{N}\mathbf{\Psi}(Au) +$$

$$+ \Big(\hat{A}^{-1}\mathbf{z}\mathbf{D} + \hat{A}^{-2}\mathbf{q}\Big)\mathbf{F}(Au) + \hat{A}^{-2}\mathbf{g}\mathbf{F}(A^{2}u) + \hat{A}^{-2}f. (23)$$

Further acting by functionals F and  $\Psi$  we get the next system

$$\begin{split} \mathbf{F}(Au) &= \mathbf{F}(\mathbf{z}) \Big[ \mathbf{DF}(Au) + \mathbf{N\Psi}(Au) \Big] + \\ &+ \mathbf{F} \Big( \hat{A}^{-1} \mathbf{q} \Big) \mathbf{F}(Au) + \mathbf{F} \Big( \hat{A}^{-1} \mathbf{g} \Big) \mathbf{F} \Big( A^{2}u \Big) + \mathbf{F} \Big( \hat{A}^{-1} f \Big), \\ \mathbf{\Psi}(u) &= \mathbf{\Psi}(\mathbf{z}) \mathbf{N\Psi}(u) + \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{z} \Big) \mathbf{N\Psi}(Au) + \\ &+ \Big[ \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{z} \Big) \mathbf{D} + \mathbf{\Psi} \Big( \hat{A}^{-2} \mathbf{q} \Big) \Big] \mathbf{F}(Au) + \\ &+ \mathbf{\Psi} \Big( \hat{A}^{-2} \mathbf{g} \Big) \mathbf{F} \Big( A^{2}u \Big) + \mathbf{\Psi} \Big( \hat{A}^{-2} f \Big), \\ \mathbf{\Psi}(Au) &= \mathbf{\Psi}(\mathbf{z}) \Big[ \mathbf{DF}(Au) + \mathbf{N\Psi}(Au) \Big] + \\ &+ \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{q} \Big) \mathbf{F}(Au) + \mathbf{\Psi} \Big( A^{-1} \mathbf{g} \Big) \mathbf{F} \Big( A^{2}u \Big) + \mathbf{\Psi} \Big( \hat{A}^{-1} f \Big), \\ \mathbf{F}(A^{2}u) &= \mathbf{F}(\mathbf{q}) \mathbf{F}(Au) + \mathbf{F}(\mathbf{g}) \mathbf{F}(A^{2}u) + \mathbf{F}(f), \text{ or } \\ &- \mathbf{F}(\mathbf{z}) \mathbf{N\Psi}(Au) + \Big[ \mathbf{I}_{n} - \mathbf{F}(\mathbf{z}) \mathbf{D} - \mathbf{F} \Big( \hat{A}^{-1} \mathbf{q} \Big) \Big] \mathbf{F}(Au) - \\ &- \mathbf{F} \Big( \hat{A}^{-1} \mathbf{z} \Big) \mathbf{D} + \mathbf{\Psi} \Big( \hat{A}^{-2} \mathbf{q} \Big) \Big] \mathbf{F}(Au) - \\ &- \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{z} \Big) \mathbf{D} + \mathbf{\Psi} \Big( \hat{A}^{-2} \mathbf{q} \Big) \Big] \mathbf{F}(Au) - \\ &- \mathbf{\Psi} \Big( \hat{A}^{-2} \mathbf{g} \Big) \mathbf{F} \Big( A^{2}u \Big) = \mathbf{\Psi} \Big( \hat{A}^{-2} f \Big), \\ \mathbf{V\Psi}(Au) - \Big[ \mathbf{\Psi}(\mathbf{z}) \mathbf{D} + \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{q} \Big) \Big] \mathbf{F}(Au) - \\ &- \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{g} \Big) \mathbf{F} \Big( A^{2}u \Big) = \mathbf{\Psi} \Big( \hat{A}^{-1} f \Big), \\ &- \mathbf{F}(\mathbf{q}) \mathbf{F}(Au) + [\mathbf{I}_{n} - \mathbf{F}(\mathbf{g})] \mathbf{F}(A^{2}u) = \mathbf{F}(f). \end{split}$$

Using the notations (19) from the above equations we get the system

$$\begin{pmatrix}
\mathbf{0}_{n} & -\mathbf{F}(\mathbf{z})\mathbf{N} & \mathbf{K}_{1} & -\mathbf{F}(\hat{A}^{-1}\mathbf{g}) \\
\mathbf{V} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{N} & -\mathbf{K}_{3} & -\mathbf{\Psi}(\hat{A}^{-2}\mathbf{g}) \\
\mathbf{0}_{n} & \mathbf{V} & -\mathbf{K}_{2} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{g}) \\
\mathbf{0}_{n} & \mathbf{0}_{n} & -\mathbf{F}(\mathbf{q}) & W
\end{pmatrix} \times \begin{pmatrix}
\mathbf{\Psi}(u) \\
\mathbf{\Psi}(Au) \\
\mathbf{F}(Au) \\
\mathbf{F}(A^{2}u)
\end{pmatrix} = \begin{pmatrix}
\mathbf{F}(\hat{A}^{-1}f) \\
\mathbf{\Psi}(\hat{A}^{-2}f) \\
\mathbf{\Psi}(\hat{A}^{-1}f) \\
\mathbf{F}(f)
\end{pmatrix}.$$
(24)

Let  $u \in \ker B_1$ . Then in the systems (23), (24) f = 0 and from (24) we get  $\operatorname{Lcol}(\Psi(u), \Psi(Au), F(Au), F(A^2u)) = 0$ , which since  $\det \mathbf{L} \neq \mathbf{0}$ , yields  $\Psi(u) = \Psi(Au) = F(Au) = F(A^2u) = 0$ . Substitution of these values into (16), (17) imply  $B_1u = A^2u = 0$ ,  $\Phi(u) = \Phi(Au) = 0$ . Taking into account (9) we acquire  $u \in \mathfrak{D}(\hat{A}^2)$  and  $B_1u = \hat{A}^2u = 0$ . By hypothesis  $\hat{A}$  is correct and so u = 0. Thus  $\ker B_1 = \{0\}$  and  $B_1$  is injective.

Conversely. Let  $\det \mathbf{L} = \mathbf{0}$ . Then there exists a vector  $\mathbf{c} = col(\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_3, \mathbf{c}_4)$ , where  $\mathbf{c}_i = col(\mathbf{c}_{i1}, ..., \mathbf{c}_{in})$ , i = 1, ..., 4 such that  $\mathbf{c} \neq \mathbf{0}$  and  $\mathbf{L}\mathbf{c} = \mathbf{0}$ , which since (24) yields

$$-\mathbf{F}(\mathbf{z})\mathbf{N}\mathbf{c}_2 + \mathbf{K}_1\mathbf{c}_3 - \mathbf{F}(\hat{A}^{-1}\mathbf{g})\mathbf{c}_4 = \mathbf{0}; \qquad (25)$$

$$\mathbf{V}\mathbf{c}_{1} - \mathbf{\Psi}\left(\hat{A}^{-1}\mathbf{z}\right)\mathbf{N}\mathbf{c}_{2} - \mathbf{K}_{3}\mathbf{c}_{3} - \mathbf{\Psi}\left(\hat{A}^{-2}\mathbf{g}\right)\mathbf{c}_{4} = \mathbf{0}; (26)$$

$$\mathbf{V}\mathbf{c}_{2} - \mathbf{K}_{2}\mathbf{c}_{3} - \mathbf{\Psi}(\hat{\mathbf{A}}^{-1}\mathbf{g})\mathbf{c}_{4} = \mathbf{0}; \tag{27}$$

$$-\mathbf{F}(\mathbf{q})\mathbf{c}_3 + \mathbf{W}\mathbf{c}_4 = \mathbf{0}. \tag{28}$$

Consider the element

$$u_0 = \mathbf{z} \mathbf{N} \mathbf{c}_1 + \hat{A}^{-1} \mathbf{z} (\mathbf{N} \mathbf{c}_2 + \mathbf{D} \mathbf{c}_3) + \hat{A}^{-2} (\mathbf{q} \mathbf{c}_3 + \mathbf{g} \mathbf{c}_4).$$
 (29)

Note that  $u_0 \neq 0$ , otherwise because of the linear independence of the vectors  $\mathbf{q}$ ,  $\mathbf{g}$ ,  $\mathbf{z}$  and  $\mathbf{D}(\hat{A}) \cap \ker A = \{0\}$  [18], we get  $\mathbf{Nc}_1 = \mathbf{Nc}_2 = \mathbf{c}_3 = \mathbf{c}_4 = \mathbf{0}$ . Then from (27) follows that  $\mathbf{c}_2 = \mathbf{0}$  and from (26) we obtain  $\mathbf{c}_1 = \mathbf{0}$ . Thus  $\mathbf{c}_i = \mathbf{0}$ , i = 1, ..., 4 and  $\mathbf{c} = \mathbf{0}$ . But the last contradicts the hypothesis  $\mathbf{c} \neq \mathbf{0}$ . So  $u_0 \neq \mathbf{0}$ . From (29), since  $\mathbf{\Phi}(\mathbf{z}) = \mathbf{I}_m$ ,  $\mathbf{K}_3 = \mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{D} + \mathbf{\Psi}(\hat{A}^{-2}\mathbf{q})$  and (26) we get

$$\begin{split} Au_0 = \mathbf{z} \big( \mathbf{N} \mathbf{c}_2 + \mathbf{D} \mathbf{c}_3 \big) + \hat{A}^{-1} \big( \mathbf{q} \mathbf{c}_3 + \mathbf{g} \mathbf{c}_4 \big), \\ A^2 u_0 = \mathbf{q} \mathbf{c}_3 + \mathbf{g} \mathbf{c}_4, \\ \Phi \big( u_0 \big) - \mathbf{N} \Psi \big( u_0 \big) = \mathbf{N} \mathbf{c}_1 - \mathbf{N} \Psi \big( \mathbf{z} \big) \mathbf{N} \mathbf{c}_1 - \mathbf{N} \Psi \Big( \hat{A}^{-1} \mathbf{z} \Big) \times \\ \times \big( \mathbf{N} \mathbf{c}_2 + \mathbf{D} \mathbf{c}_3 \big) - \mathbf{N} \Psi \Big( \hat{A}^{-2} \mathbf{q} \big) \mathbf{c}_3 - \mathbf{N} \Psi \Big( \hat{A}^{-2} \mathbf{g} \big) \mathbf{c}_4 = \end{split}$$

$$= \mathbf{N} \left[ \mathbf{V} \mathbf{c}_1 - \mathbf{\Psi} (\hat{A}^{-1} \mathbf{z}) \mathbf{N} \mathbf{c}_2 - \mathbf{K}_3 \mathbf{c}_3 - \mathbf{\Psi} (\hat{A}^{-2} \mathbf{g}) \mathbf{c}_4 \right] =$$

$$= \mathbf{N} \mathbf{0} = \mathbf{0}.$$

Then  $\Phi(u_0) = \mathbf{N}\Psi(u_0)$  and so  $u_0$  satisfies the first boundary condition (17). We will show, using (27) and (25), that  $u_0$  satisfies the second boundary condition (17)

$$\begin{split} & \Phi \big(Au_0 \, \big) - \mathbf{DF} \big(Au_0 \, \big) - \mathbf{N\Psi} \big(Au_0 \, \big) = \mathbf{Nc}_2 + \mathbf{Dc}_3 \, - \\ & - \mathbf{DF} \big(\mathbf{z} \, \big) \big( \mathbf{Nc}_2 + \mathbf{Dc}_3 \big) - \mathbf{DF} \Big( \hat{A}^{-1} \mathbf{q} \, \Big) \mathbf{c}_3 - \mathbf{DF} \Big( \hat{A}^{-1} \mathbf{g} \, \Big) \mathbf{c}_4 \, - \\ & - \mathbf{N\Psi} \big(\mathbf{z} \, \big) \big( \mathbf{Nc}_2 + \mathbf{Dc}_3 \, \big) - \mathbf{N\Psi} \Big( \hat{A}^{-1} \mathbf{q} \, \Big) \mathbf{c}_3 - \mathbf{N\Psi} \Big( \hat{A}^{-1} \mathbf{g} \, \Big) \mathbf{c}_4 = \\ & = \mathbf{N} \bigg[ \mathbf{V} \mathbf{c}_2 - \mathbf{K}_2 \mathbf{c}_3 - \mathbf{\Psi} \Big( \hat{A}^{-1} \mathbf{g} \, \Big) \mathbf{c}_4 \, \bigg] + \\ & + \mathbf{D} \bigg[ - \mathbf{F} \big(\mathbf{z} \, \big) \mathbf{Nc}_2 + \mathbf{K}_1 \mathbf{c}_3 - \mathbf{F} \Big( \hat{A}^{-1} \mathbf{g} \, \Big) \mathbf{c}_4 \, \bigg] = \mathbf{NO} + \mathbf{DO} = \mathbf{0}, \end{split}$$

where  $\mathbf{K}_1$ ,  $\mathbf{K}_2$  from (19). So  $u_0\in \mathfrak{D}(B_1)$ . Now, using (25) and (28) we will show that  $u_0\in \ker B_1$ 

$$\begin{split} B_1 u_0 &= A^2 u_0 - \mathbf{q} \mathbf{F} \big( A u_0 \big) - \mathbf{g} \mathbf{F} \Big( A^2 u_0 \Big) = \mathbf{q} \mathbf{c}_3 + \mathbf{g} \mathbf{c}_4 - \\ &- \mathbf{q} \Big[ \mathbf{F} \big( \mathbf{z} \big) \big( \mathbf{N} \mathbf{c}_2 + \mathbf{D} \mathbf{c}_3 \big) + \mathbf{F} \Big( \hat{A}^{-1} \mathbf{q} \Big) \mathbf{c}_3 + \mathbf{F} \Big( \hat{A}^{-1} \mathbf{g} \Big) \mathbf{c}_4 \Big] - \\ &- \mathbf{g} \mathbf{F} \big( \mathbf{q} \big) \mathbf{c}_3 - \mathbf{g} \mathbf{F} \big( \mathbf{g} \big) \mathbf{c}_4 = \\ &= \mathbf{q} \Big[ - \mathbf{F} \big( \mathbf{z} \big) \mathbf{N} \mathbf{c}_2 + \mathbf{K}_1 \mathbf{c}_3 - \mathbf{F} \Big( \hat{A}^{-1} \mathbf{g} \Big) \mathbf{c}_4 \Big] + \\ &+ \mathbf{g} \Big[ - \mathbf{F} \big( \mathbf{q} \big) \mathbf{c}_3 + \mathbf{W} \mathbf{c}_4 \Big] = \mathbf{q} \mathbf{0} + \mathbf{g} \mathbf{0} = \mathbf{0}. \end{split}$$

So there exists a nonzero element  $u_0 \in \mathfrak{D}(B_1)$  and  $u_0 \in \ker B_1$ . This means that  $B_1$  is not injective. Hence the operator  $B_1$  is injective if and only if  $\det \mathbf{L} \neq \mathbf{0}$ .

(ii) Since  $\det \mathbf{L} \neq \mathbf{0}$ , the system (24) for all  $f \in X$  has an unique solution

$$col(\Psi(u), \Psi(Au), \mathbf{F}(Au), \mathbf{F}(A^{2}u)) =$$

$$= \mathbf{L}^{-1}col(\mathbf{F}(\hat{A}^{-1}f), \Psi(\hat{A}^{-2}f), \Psi(\hat{A}^{-1}f), \mathbf{F}(f)) \quad (30)$$

and the operator  $B_1$ , by statement (i), is injective. Substituting (30) into (23) we obtain the unique solution (20) of the problem (16), (17). In the above solution an element f is arbitrary. Consequently,  $R(B_1) = X$ . Since the operators  $\hat{A}^{-2}$ ,  $\hat{A}^{-1}$  and the functional vectors  $\mathbf{F}$  and  $\Psi$  are bounded, from (20) follows the boundedness of  $B_1^{-1}$ , i. e. the operator  $B_1$  is correct. The theorem is proved.

The next corollary follows from the above theorem for  $\mathbf{q}=\mathbf{g}=\mathbf{0}$  and is useful for solving some classes of differential equations with nonlocal boundary conditions.

**Corollary 2.** Let the operators A,  $\hat{A}$ , the vectors  $\mathbf{z}$ ,  $\Phi$ ,  $\Psi$ ,  $\mathbf{F}$ ,  $\mathbf{V}$  and matrices  $\mathbf{D}$ ,  $\mathbf{N}$  be defined as in Theorem 2 and the operator  $B_1: X \to X$  be defined by

$$B_1 u = A^2 u = f,$$

$$\mathfrak{D}(B_1) = \{ u \in \mathfrak{D}(A^2) : \Phi(u) = \mathbf{N}\Psi(u), \qquad (31)$$

$$\Phi(Au) = \mathbf{D}\mathbf{F}(Au) + \mathbf{N}\Psi(Au) \}.$$

Then

(i) The operator  $B_1$  corresponding to the problem (31) is injective if and only if

$$\det \mathbf{L}_{1} = \det \begin{pmatrix} \mathbf{0}_{n} & -\mathbf{F}(\mathbf{z})\mathbf{N} & \mathbf{I}_{n} - \mathbf{F}(\mathbf{z})\mathbf{D} \\ \mathbf{V} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{N} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{D} \\ \mathbf{0}_{n} & \mathbf{V} & -\mathbf{\Psi}(\mathbf{z})\mathbf{D} \end{pmatrix} \neq \mathbf{0}. (32)$$

(ii) If the operator  $B_1$  is injective, then it is correct and the unique solution of (31) is given by

$$u = B_1^{-1} f = \hat{A}^{-2} f + \left( \mathbf{z} \mathbf{N}, \, \hat{A}^{-1} \mathbf{z} \mathbf{N}, \, \hat{A}^{-1} \mathbf{z} \mathbf{D} \right) \times$$
$$\times \mathbf{L}_1^{-1} col\left( \mathbf{F} \left( \hat{A}^{-1} f \right), \, \mathbf{\Psi} \left( \hat{A}^{-2} f \right), \, \mathbf{\Psi} \left( \hat{A}^{-1} f \right) \right). \tag{33}$$

*Proof*: (i) For g = q = 0 from (18) and (19) immediately follows

$$\det \mathbf{L} = \det \begin{pmatrix} \mathbf{0}_{n} & -\mathbf{F}(\mathbf{z})\mathbf{N} & \mathbf{I}_{n} - \mathbf{F}(\mathbf{z})\mathbf{D} & \mathbf{0}_{n} \\ \mathbf{V} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{N} & -\mathbf{\Psi}(\hat{A}^{-1}\mathbf{z})\mathbf{D} & \mathbf{0}_{n} \\ \mathbf{0}_{n} & \mathbf{V} & -\mathbf{\Psi}(\mathbf{z})\mathbf{D} & \mathbf{0}_{n} \\ \mathbf{0}_{n} & \mathbf{0}_{n} & \mathbf{0}_{n} & \mathbf{I}_{n} \end{pmatrix} . (34)$$

It is evident that  $\det L = \det L_1$ . From (20) for g = q = 0 follows the solution of (31)

$$u = B_1^{-1} f = \hat{A}^{-2} f + \left( \mathbf{zN}, \, \hat{A}^{-1} \mathbf{zN}, \, \hat{A}^{-1} \mathbf{zD}, \, 0 \right) \times$$
$$\times \mathbf{L}^{-1} col\left( \mathbf{F} \left( \hat{A}^{-1} f \right), \, \Psi \left( \hat{A}^{-2} f \right), \, \Psi \left( \hat{A}^{-1} f \right), \, \mathbf{F} (f) \right). \tag{35}$$

It is easy to verify that

$$\begin{split} &\left(\mathbf{z}\mathbf{N},\ \hat{A}^{-1}\mathbf{z}\mathbf{N},\,\hat{A}^{-1}\mathbf{z}\mathbf{D},\ 0\right)\times\\ &\times\mathbf{L}^{-1}col\Big(\mathbf{F}\Big(\hat{A}^{-1}f\Big),\,\mathbf{\Psi}\Big(\hat{A}^{-2}f\Big),\,\mathbf{\Psi}\Big(\hat{A}^{-1}f\Big),\,\mathbf{F}(f)\Big) =\\ &=\Big(\mathbf{z}\mathbf{N},\,\hat{A}^{-1}\mathbf{z}\mathbf{N},\,\hat{A}^{-1}\mathbf{z}\mathbf{D}\Big)\times\\ &\times\mathbf{L}_{1}^{-1}col\Big(\mathbf{F}\Big(\hat{A}^{-1}f\Big),\,\mathbf{\Psi}\Big(\hat{A}^{-2}f\Big),\,\mathbf{\Psi}\Big(\hat{A}^{-1}f\Big)\Big). \end{split}$$

Hence, from (35) follows (33).

# Examples

In the next example we use the extension method from Theorem 1.

**Example.** The multipoint problem for loaded integro-differential equation on  $\mathbb{C}[0, 1]$ 

# ΤΕΟΡΕΤИЧЕСКАЯ И ΠΡИКΛΑΔΗΑЯ ΜΑΤΕΜΑΤИΚΑ

$$u'' - 3t \int_{0}^{1} x^{2} u''(x) dx + \frac{1}{2} (t^{2} + 1) [u'(1) - u'(0)] =$$

$$= 8t^{2} + 2t + 12,$$

$$u(0) = \frac{1}{6} u(1/2) + \frac{1}{18} u(1), \ u'(0) = \frac{2}{9} u(1)$$
(36)

is correct and the unique solution of (36) is given by

$$u(t) = 4t^3 + 2t^2 + 2t + 1. (37)$$

*Proof*: If we compare (36) with (10), it is natural to take Au = u''(t),  $\mathfrak{D}(A) = \{u \in \mathbb{C}^2[0, 1]\}$ ,  $X_A^2 = \mathbb{C}^2[0, 1]$ ,  $X_A^1 = \mathbb{C}^1[0, 1]$ , m = n = 2,  $\mathbf{z} = (z_1, z_2) = (1, t)$ ,  $\hat{A}u = Au$ ,

$$\mathfrak{D}(\hat{A}) = \left\{ u \in \mathfrak{D}(A) : u(0) = u'(0) = 0 \right\},$$

$$Bu = u'' - 3t \int_{0}^{1} x^{2} u''(x) dx + \frac{1}{2} (t^{2} + 1) [u'(1) - u'(0)] =$$

$$= u'' - 3t \int_{0}^{1} x^{2} u''(x) dx + \frac{1}{2} (t^{2} + 1) \int_{0}^{1} u''(x) dx,$$

$$\mathfrak{D}(B) =$$

$$= \left\{ u(x) \in \mathfrak{D}(A) : \begin{pmatrix} u(0) \\ u'(0) \end{pmatrix} = \begin{pmatrix} 1/6 & 1/18 \\ 0 & 2/9 \end{pmatrix} \begin{pmatrix} u(1/2) \\ u(1) \end{pmatrix} \right\}. (38)$$

Since (5), the operator  $\hat{A}$ , is correct and its solution is  $\hat{A}^{-1}f(t)=\int_0^t (t-x)f(x)\mathrm{d}x$ . Further comparing (36), (38) with (10), we take  $g_1=3t$ ,  $g_2=-\frac{1}{2}(t^2+1)$ ,  $f=8t^2+2t+12$ ,

$$\mathbf{N} = \begin{pmatrix} 1/6 & 1/18 \\ 0 & 2/9 \end{pmatrix}, \quad F_1(Au) = \int_0^1 x^2 u''(x) \, \mathrm{d}x,$$
$$F_2(Au) = \int_0^1 u''(x) \, \mathrm{d}x.$$

Then

$$F_1(f) = \int_0^1 x^2 f(x) dx, \quad F_2(f) = \int_0^1 f(x) dx,$$

$$\Phi(u) = \begin{pmatrix} \Phi_1(u) \\ \Phi_2(u) \end{pmatrix} = \begin{pmatrix} u(0) \\ u'(0) \end{pmatrix},$$

$$\Psi(u) = \begin{pmatrix} \Psi_1(u) \\ \Psi_2(u) \end{pmatrix} = \begin{pmatrix} u(1/2) \\ u(1) \end{pmatrix}.$$

The set  $\mathbf{z}=(1,t)$  is biorthogonal to  $(\Phi_1,\Phi_2)$ . From  $|\Psi_1(u)|=|u(1/2)|\leq \|u\|_{\mathbb{C}}+\|u'\|_{\mathbb{C}}=\|u\|_{\mathbb{C}^1}$  follows that  $\Psi_1\in\mathbb{C}^{1^*}=X_A^{m-1^*}=X_A^{1^*}$ . By analogy  $\Psi_2,\ \Psi_i\in\mathbb{C}^{1^*},\ i=1,\ 2.$  Further from  $\left|F_1(f)\right|=\left|\int_0^1x^2f(x)\mathrm{d}\,x\right|\leq \|f\|_{\mathfrak{C}}$  it follows that  $F_1\in\mathbb{C}\left[0,1\right]^*=X^*$ . By analogy it is

proved that  $F_2 \in X^*$ . We can apply Theorem 1. Now we calculate

$$\hat{A}^{-1}g_1(t) = \int_0^t (t-x)g_1(x) dx = \int_0^t (t-x)3x dx = \frac{t^3}{2},$$

$$\hat{A}^{-1}g_2(t) = -\frac{1}{2} \int_0^t (t-x)(x^2+1) dx = -\frac{t^2(t^2+6)}{24}.$$

Compute

$$\hat{A}^{-1}g = (\hat{A}^{-1}g_1, \hat{A}^{-1}g_2) = \left(\frac{t^3}{2}, -\frac{t^2(t^2+6)}{24}\right).$$

Further we find  $\Psi_1(z_1)=z_1(1/2)=1$ ,  $\Psi_1(z_2)=z_2(1/2)=1/2$ ,  $\Psi_2(z_1)=z_1(1)=1$ ,  $\Psi_2(z_2)=z_2(1)=1$ . Then  $\Psi(\mathbf{z})=\begin{pmatrix} 1 & 1/2 \\ 1 & 1 \end{pmatrix}$ . Further compute  $\Psi_1\left(\hat{A}^{-1}g_1\right)=1/2$ ,  $\Psi_1\left(\hat{A}^{-1}g_2\right)=-25/384$ ,  $\Psi_2\left(\hat{A}^{-1}g_2\right)=-7/24$ , then  $\Psi(\hat{A}^{-1}g)=\begin{pmatrix} 1/16 & -25/384 \\ 1/2 & -7/24 \end{pmatrix}.$ 

Now we find

$$\begin{split} F_1\left(g_1\right) &= \int_0^1 3x^3 \, \mathrm{d}\, x = 3/4, \\ F_1\left(g_2\right) &= \frac{1}{2} \int_0^1 x^2 \left(x^2 + 1\right) \! \mathrm{d}\, x = -4/15, \\ F_2\left(g_1\right) &= \int_0^1 3x \, \mathrm{d}\, x = 3/2, \\ F_2\left(g_2\right) &= -\frac{1}{2} \int_0^1 \left(x^2 + 1\right) \! \mathrm{d}\, x = -2/3. \end{split}$$

Then

$$\mathbf{F}(g) = \begin{pmatrix} 3/4 & -4/15 \\ 3/2 & -2/3 \end{pmatrix}.$$

Since

$$\mathbf{W} = \mathbf{I}_2 - \mathbf{F}(\mathbf{g}) = \begin{pmatrix} 1/4 & 4/15 \\ -3/2 & 5/3 \end{pmatrix},$$

$$\mathbf{V} = \mathbf{I}_2 - \mathbf{\Psi}(\mathbf{z})\mathbf{N} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \begin{pmatrix} 1 & 1/2 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1/6 & 1/18 \\ 0 & 2/9 \end{pmatrix} =$$

$$= \begin{pmatrix} 5/6 & -1/6 \\ -1/6 & 13/18 \end{pmatrix},$$

and  $\det \mathbf{W} \neq \mathbf{0}$ ,  $\det \mathbf{V} \neq \mathbf{0}$ , the problem (36), by Theorem 1 (ii), is correct. For  $f = 8t^2 + 2t + 12$  we calculate

$$\hat{A}^{-1}f(t) = \int_0^t (t-x)(2x+4) dx = \frac{t^2(2t^2+t+18)}{3},$$

$$F_1(f) = \int_0^1 x^2 (8x^2+2x+12) dx = 61/10,$$

$$F_2(f) = \int_0^1 (8x^2+2x+12) dx = 47/3.$$

Then  $\mathbf{F}(f) = col(61/10, 47/3)$ . We also compute

$$\Psi_1(\hat{A}^{-1}f) = \hat{A}^{-1}f_{|t=1/2} = 19/12,$$
  
$$\Psi_2(\hat{A}^{-1}f) = \hat{A}^{-1}f_{|t=1} = 7.$$

Then

$$\Psi(\hat{A}^{-1}f) = col(19/12,7).$$

Substitution of these values into (12) yields the solution to the problem (36)

$$u(t) = \hat{A}^{-1}f + [\hat{A}^{-1}g + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\Psi(A^{-1}g)]\mathbf{W}^{-1}\mathbf{F}(f) + \mathbf{z}\mathbf{N}\mathbf{V}^{-1}\Psi(\hat{A}^{-1}f) = \frac{t^2(2t^2 + t + 18)}{3} + \mathbf{v}^{-1}\Psi(\hat{A}^{-1}f) = \frac{t^2($$

### References

- Bloom F. Ill posed Problems for Integrodifferential Equations in Mechanics and Electromagnetic Theory. SIAM, 1981. 232 p.
- Cushing J. M. Integrodifferential Equations and Delay Models in Population Dynamics. Springer, 1977. 198 p.
- Apreutesei N., Ducrot A., Volpert V. Travelling waves for integro-differential equations in population dinamics. *Discrete Cont. Dyn. Syst.*, Ser. B 11, 2009, no. 3, pp. 541–561.
- Arisawa M. A remark on the definitions of viscosity solutions for the integro-differential equations with Lévy operators. J. Math. Pures Appl., 2008, vol. 89, pp. 567-574.
- 5. Cannon J. R. The solution of the heat equation subject to the specification of energy. *Quart. Appl. Math.*, 1963, vol. 21, pp. 155–160.
- Ionkin N. I. Solution of one boundary value problem of heat conduction theory with a nonclassical boundary condition. *Differencial'nye uravneniya* [Differential Equations], 1977, vol. 13, no. 2, pp. 294–304 [In Russian].
- Kamynin L. I. On a boundary problem in the theory of heat conduction with a nonclassical boundary condition. ZHurnal vychislitel'noj matematiki i matematicheskoj fiziki [Computational Mathematics and Mathematical Physics], 1964, vol. 4, no. 6, pp. 1006– 1024 [In Russian].

$$\begin{split} &+\left[\left(\frac{t^3}{2},-\frac{t^2\Big(t^2+6\Big)}{24}\right) + \left(1,t\right) \left(\frac{1/6}{0} \ \frac{1/18}{2/9}\right) \times \right. \\ &\times \left(\frac{5/6}{-1/6} \ \frac{-1/6}{13/18}\right)^{-1} \left(\frac{1/16}{1/2} \ \frac{-25/384}{-7/24}\right) \right] \frac{1}{49} \left(\frac{100}{90} \ \frac{-16}{15}\right) \times \\ &\times \left(\frac{61/10}{47/3}\right) + \left(1,t\right) \left(\frac{1/6}{0} \ \frac{1/18}{2/9}\right) \left(\frac{5/6}{-1/6} \ \frac{-1/6}{13/18}\right)^{-1} \times \\ &\times \left(\frac{19/12}{7}\right) = 4t^3 + 2t^2 + 2t + 1. \end{split}$$

#### Conclusion

The main results of this paper are Theorems 1 and 2, where the problems Bu = f,  $B_1u = f$  are solved by extension method. This method is essentially simpler and more convenient in the case of quadratic operator  $B_1 = B^2$ . In this case the solvability condition and a solution of  $B_1u = f$  can be obtained by application of the formula for solution of Bu = f twice. The upcoming Part 2 of this paper will be devoted to decomposition method for this case. Note that the extension method is a generalization of direct method which is presented in [30]. The essential ingredient in our approach is the extension of the main idea in [26].

- 8. Kandemir M. Nonlocal boundary value problems with transmission conditions. *Gulf Journal of Mathematics*, 2015, vol. 3, iss. 1, no. 3, pp. 1–17.
- 9. Medlock J., Kot M. Spreading disease: integro-differential equations old and new. *Mathematical Biosciences*, Elsevier, 2003, vol. 184, pp. 201–222.
- 10. Samarskii A. A. On certain problems of the modern theory of differential equations. *Differencial'nye* uravneniya [Differential Equations], 1980, vol. 16, no. 11, pp. 1221–1228 [In Russian].
- 11. Sachs E. W., Strauss A. K. Efficient solution of a partial integro-differential equation in finance. *Appl. Numer. Math.*, 2008, vol. 58, pp. 1687–1703.
- 12. Schumacher K. Traveling front solutions for integro-differential equations. I. J. Reine Angew. Math., 1980, vol. 316, pp. 54–70.
- 13. Shivanian E. Analysis of meshless local radial point interpolation (MLRPI) on a nonlinear partial integro-differential equation arising in population dynamics. *Engineering Analysis with Boundary Elements*, 2003, vol. 37, pp. 1693–1702.
- 14. Tersenov Alkis S. Ultraparabolic equations and unsteady heat transfer. *Journal of Evaluation Equations*, 2005, vol. 5, no. 2, pp. 277–289.
- 15. Tamarkin J. D. The notion of the Green's function in the theory of integro-differential equations. *Trans. Amer. Math. Soc.*, 1927, vol. 29, pp. 755–800.
- 16. Bitsatze A. V., Samarskii A. A. On some simplest generalization of linear elliptic problems. *Doklady AN*

### ΤΕΟΡΕΤИЧЕСКАЯ И ΠΡИΚΛΑΔΗΑЯ ΜΑΤΕΜΑΤИΚΑ

- SSSR [Reports of the USSR Academy of Sciences], 1969, vol. 185, pp. 739–740 [In Russian].
- 17. Il'in V. A., Moiseev E. L. Two dimensial nonlocal boundary value problem for Poissons operator in differential and difference variants. *Matematicheskoe modelirovanie* [Mathematical Models and Computer Simulations], 1990, vol. 2, no. 8, pp. 132–156 [In Russian].
- 18. Avalishvili G., Avalishvili M., Gordeziani D. On a nonlocal problem with integral boundary conditions for a multidimensional elliptic equation. *Applied Mathemat. Letters*, 2004, vol. 24, no. 4, pp. 566-571.
- 19. Kalmenov T. S., Tokmaganbetov N. E. On a nonlocal boundary value problem for the multidimensioal heat equation in a noncylindrical domain. S. M. J., 2013, vol. 54, no. 6, pp. 1287–1293.
- 20. Sadybekov M. A., Turmetov B. K. On an analog of periodic boundary value problems for the Poisson equation in the disk. *Differencial'nye uravneniya* [Differential Equations], 2014, vol. 50, pp. 264–268 [In Russian].
- 21. Pulkina L. S. A nonlocal problem with integral condition for a hyperbolic equation. *Differencial'nye uravneniya* [Differential Equations], 2004, vol. 40, no. 7, pp. 15–27 [In Russian].
- 22. Abdullaev A. R., Skachkova E. A. On one class of multipoint boundary value problems for a second-order linear functional-differential equation. *Journal of Mathematical Sciences*, 2018, vol. 230, no. 5, pp. 647–650.
- 23. Benchohra M., Ntouyas S. K. Existence results on the semiinfinite interval for first and second order integrodifferential equations in banach spaces with non-local conditions. *Acta Univ. Palacki. Olomuc, Fac. rer. nat. V Mathematica*, 2002, vol. 41, pp. 13–19.
- 24. Ntouyas S. A six-point boundary value problem of nonlinear coupled sequential fractional integro-differential equations and coupled integral boundary conditions. *Journal of Applied Mathematics and Computing*, 2018, vol. 56, no. 1-2, pp. 367–389.
- 25. Georgiou D. N., Kougias I. E. On fuzzy fredholm and voltera integral equations. *Journal of Fuzzy Mathematics*, 2001, vol. 9, no. 4, pp. 943–951.

- 26. Oinarov R. O., Parasidi I. N. Correct extensions of operators with finite defect in Banach spases. *Izvestiya Akademii nauk Kazahskoj SSR*, 1988, vol. 5, pp. 42–46 [In Russian].
- 27. Parasidis I. N. and Tsekrekos P. C. Correct and self-adjoint problems for quadratic operators. *Eurasian Mathematical Journal*, 2010, vol. 1, no. 2. pp. 122–135.
- 28. Parasidis I. N., Providas E. Extension operator method for the exact solution of integro-differential equations. In: Pardalos P., Rassias T. (eds). Contributions in Mathematics and Engineering: In Honor of Constantin Caratheodory. Springer, Cham., 2016, pp. 473–496.
- 29. Polyanin A. D., Zhurov A. I. Exact solutions to some classes of nonlinear integral, integro-functional and integro-differential equation. *Dokl. Math.*, 2008, vol. 77, no. 2, pp. 315–319.
- Wazwaz A. M. Linear and Nonlinear Integral Equations: Methods and Applications. Springer, Beijing, 2011. 657 p.
- 31. Dezin A. A. Nonstandard problems. *Matematicheskie zametki*, 1987, vol. 41, no. 3, pp. 356–364 [In Russian].
- 32. Krein M. G. The theory of self-adjoint extensions of semi-bounded Hermitian operators and its aplications. *Matematicheskij sbornik*, 1947, vol. 20, no. 3, pp. 431–495 [In Russian].
- 33. Neumann J. Von. Allgemeine eigenwerttheorie hermitescher functional operatoren. *Math. Ann.*, Bd., 1929–1930, vol. 102, pp. 49–131.
- 34. Kokebaev B. K., Otelbaev M., Shynybekov A. N. About restrictions and extensions of operators. *Doklady AN SSSR* [Reports of the USSR Academy of Sciences], 1983, vol. 271, no. 6, pp. 1307–1310 [In Russian].
- 35. Parassidis I. N. and Tsekrekos P. C. Correct selfadjoint and positive extensions of nondensely defined symmetric operators. *Abstract and Applied Analysis*, 2005, no. 7, pp. 767–790.

УДК 338.984

doi:10.31799/1684-8853-2018-6-14-23

Метод нахождения точных решений для интегро-дифференциальных уравнений Фредгольма с многоточечными и интегральными краевыми условиями. Часть 1. Метод расширения

- H. H. Васильев $^{a,6}$ , канд. физ.-мат. наук, старший научный сотрудник, orcid.org/0000-0002-0841-1168, vasiliev@pdmi.ras.ru
- И. Н. Парасидис<sup>в</sup>, PhD, доцент, paras@teilar.gr
- E. Провидас<sup>г</sup>, PhD, доцент, providas@teilar.gr
- <sup>а</sup>Санкт-Петербургское отделение Математического института им. В. А. Стеклова РАН, наб. р. Фонтанки, 27, Санкт-Петербург, 191023, РФ
- <sup>6</sup>Санкт-Петербургский государственный электротехнический университет «ЛЭТИ», Санкт-Петербург, ул. Профессора Попова, 5, Санкт-Петербург, 197376, РФ
- вКафедра электротехники, Технологический институт Фессалии, 41110, Лариса, Греция
- <sup>г</sup>Кафедра машиностроения, Технологический институт Фессалии, 41110, Лариса, Греция

Введение: краевые задачи для дифференциальных и интегро-дифференциальных уравнений с многоточечными и нелокальными граничными условиями возникают в различных областях механики, физики, биологии, биотехнологии, химической инже-

# ΤΕΟΡΕΤИЧЕСКАЯ И ΠΡИΚΛΑΔΗΑЯ ΜΑΤΕΜΑΤИΚΑ

нерии, медицинской науки, финансов и других. Нахождение точных решений краевых задач с фредгольмовыми интегро-дифференциальными уравнениями является трудной проблемой. В большинстве случаев решения получаются численными методами. Цель: поиск необходимых и достаточных условий разрешимости абстрактных операторных уравнений и метод построения их точных решений. Результаты: предложен прямой метод для точного решения некоторого класса обыкновенных дифференциальных или фредгольмовых интегро-дифференциальных уравнений с сепарабельными ядрами и многоточечными и интегральными граничными условиями. Исследованы абстрактные уравнения вида  $Bu = Au - g\mathbf{F}(Au) = f$  и  $B_1u = A^2u - q\mathbf{F}(Au) - g\mathbf{F}(A^2u) = f$  с нелокальными граничными условиями  $\Phi(u) = \mathbf{N}\Psi(u)$  и  $\Phi(u) = \mathbf{N}\Psi(u)$ ,  $\Phi(Au) = \mathbf{D}\mathbf{F}(Au) + \mathbf{N}\Psi(Au)$  соответственно, где  $\mathbf{q}$ ,  $\mathbf{g}$  являются векторами,  $\mathbf{D}$ ,  $\mathbf{N}$  — матрицами, а  $\mathbf{F}$ ,  $\mathbf{\Phi}$ ,  $\mathbf{\Psi}$  — функциональными векторами. Предложенный метод прост в использовании и может быть легко интегрирован в любую систему компьютерной алгебры. Исследована корректность уравнений вида Bu = f и  $B_1u = f$  и их точные решения. Вторая часть этой статьи будет посвящена случаю, когда оператор  $B_1$  имеет квадратичную факторизацию. Ключевые слова — дифференциальные и фредгольмовы интегро-дифференциальные уравнения, многоточечные и нелокаль

**Цитирование:** Vassiliev N. N., Parasidis I. N., Providas E. Exact solution method for Fredholm integro-differential equations with multipoint and integral boundary conditions. Part 1. Extension method. *Информационно-управляющие системы*, 2018, № 6, с. 14—23. doi:10.31799/1684-8853-2018-6-14-23

ные интегральные граничные условия, разложение операторов, корректность операторов, точные решения.

Citation: Vassiliev N. N., Parasidis I. N., Providas E. Exact solution method for Fredholm integro-differential equations with multipoint and integral boundary conditions. Part 1. Extension method. *Informatsionno-upravliaiushchie sistemy* [Information and Control Systems], 2018, no. 6, pp. 14–23. doi:10.31799/1684-8853-2018-6-14-23

# УВАЖАЕМЫЕ АВТОРЫ!

Научные базы данных, включая SCOPUS и Web of Science, обрабатывают данные автоматически. С одной стороны, это ускоряет процесс обработки данных, с другой — различия в транслитерации  $\Phi$ ИО, неточные данные о месте работы, области научного знания и т. д. приводят к тому, что в базах оказывается несколько авторских страниц для одного и того же человека. В результате для всех по отдельности считаются индексы цитирования, снижая рейтинг ученого.

Для идентификации авторов в сетях Thomson Reuters проводит регистрацию с присвоением уникального индекса (ID) для каждого из авторов научных публикаций.

Процедура получения ID бесплатна и очень проста, есть возможность провести регистрацию на 12-ти языках, включая русский (чтобы выбрать язык, кликните на зеленое поле вверху справа на стартовой странице): https://orcid.org